

LECTURES ABOUT (ADVANCED) STATISTICAL PHYSICS

T.S.Biró, MTA Wigner Research Centre for Physics, Budapest

Lectures given at: University of Johannesburg, South-Africa,

November 26 – November 29, 2012.

- 1. Ancient Thermodynamics (... - 1870)**
- 2. The Rise of Statistical Physics (1890 – 1920)**
- 3. Modern (postwar) Problems (1940 – 1980)**
- 4. Corrections (1950 – 2005)**
- 5. Generalizations (1960 – 2010)**
- 6. High Energy Physics (1950 – 2010)**

LECTURE TWO ABOUT (ADVANCED) STATISTICAL PHYSICS

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November 27, 2012.

Kinetic theory

- **Sum of random forces: noise (Gaussian)**
- **Brownian motion** TSB, CG, PRL 79, 3138, 1997
- **Langevin and Fokker-Planck equations**
- **Fluctuation-dissipation theorem**
- **Boltzmann equation**
- **Entropy → equilibrium theory** TSB, PG, hep-ph/0503204

General Langevin problem

$$\dot{p} = F(p, z)$$

Wang + Uehlenbeck:

use $R(p)$ test function!

Many $p(t)$ evolutions from $p(0)$: $f(p,t)$ distribution

$$\int dp R(p) f(p, t + dt) = \int dp \langle R(p + dt F(p, z)) \rangle f(p, t)$$

average over noise $\langle F \rangle = -G(p)$,

$$\langle FF \rangle - \langle F \rangle \langle F \rangle = 2 D(p) / dt$$

General Langevin problem

Expansion til $o(dt)$ gives:

$$\int dp R(p) \frac{\partial f}{\partial t} = \int dp [-R'(p)G(p) + R''(p)D(p)] f(p,t)$$

Fokker-Planck equation after partial integration:

$$\frac{\partial f}{\partial t} = \frac{\partial}{\partial p} (G(p) f) + \frac{\partial^2}{\partial p^2} (D(p) f)$$

Particular Langevin problem

TSB, G Gy, AJ, GP, JPG31, 759, 2005

$$\dot{p} = z - G(E) \frac{\partial E}{\partial p}$$

$$\langle z(t) \rangle = 0$$

$$\langle z(t)z(t') \rangle = 2 D(E) \delta(t-t')$$

In the Fokker – Planck equation:

$$D(p) = D(E)$$

$$G(p) = -G(E) \frac{\partial E}{\partial p}$$

Stationary distribution:

$$f(p) = \frac{A}{D(E)} \exp \left(- \int G(E) \frac{dE}{D(E)} \right) = A \exp \left(- \int \frac{dE}{T(E)} \right)$$

Inverse logarithmic slope \Rightarrow *temperature*

$$\frac{1}{T(E)} = -\frac{d}{dE} \ln f(E)$$

$$T(E) = \frac{D(E)}{G(E) + D'(E)}$$

$$T_{Einstein} = D(E) / G(E)$$

$$T_{Gibbs} = D(0) / G(0)$$

General inverse slope

Stationary distribution:

$$f(p) = A \exp\left(-\int \frac{dE}{T(E)}\right)$$

1) Gibbs: $T(E) = T \rightarrow \exp(-E/T)$

2) Tsallis: $T(E) = T/q + (1-1/q) E \rightarrow$
 $(1 + (q-1) E / T)^{-q / (q-1)}$

$T(T) = T$: a fixed point of the sliding slope

Fluctuation



Dissipation *theorem*

$$D_{ij}(E) = T(E) \left(G_{ij}(E) + D'_{ij}(E) \right)$$

(Hamiltonian eom does not change energy E!)

$$\dot{p}_i = -G_{ij} \nabla_j E + z_i$$

$$D_{ij}(E) = \frac{1}{f(E)} \int_E^{\infty} G_{ij}(x) f(x) dx$$

with $f(E)$ stationary distribution

Fluctuation



Dissipation *theorem*

particular cases (for constant G):

Gibbs:

$$D_{ij} = T G_{ij}$$

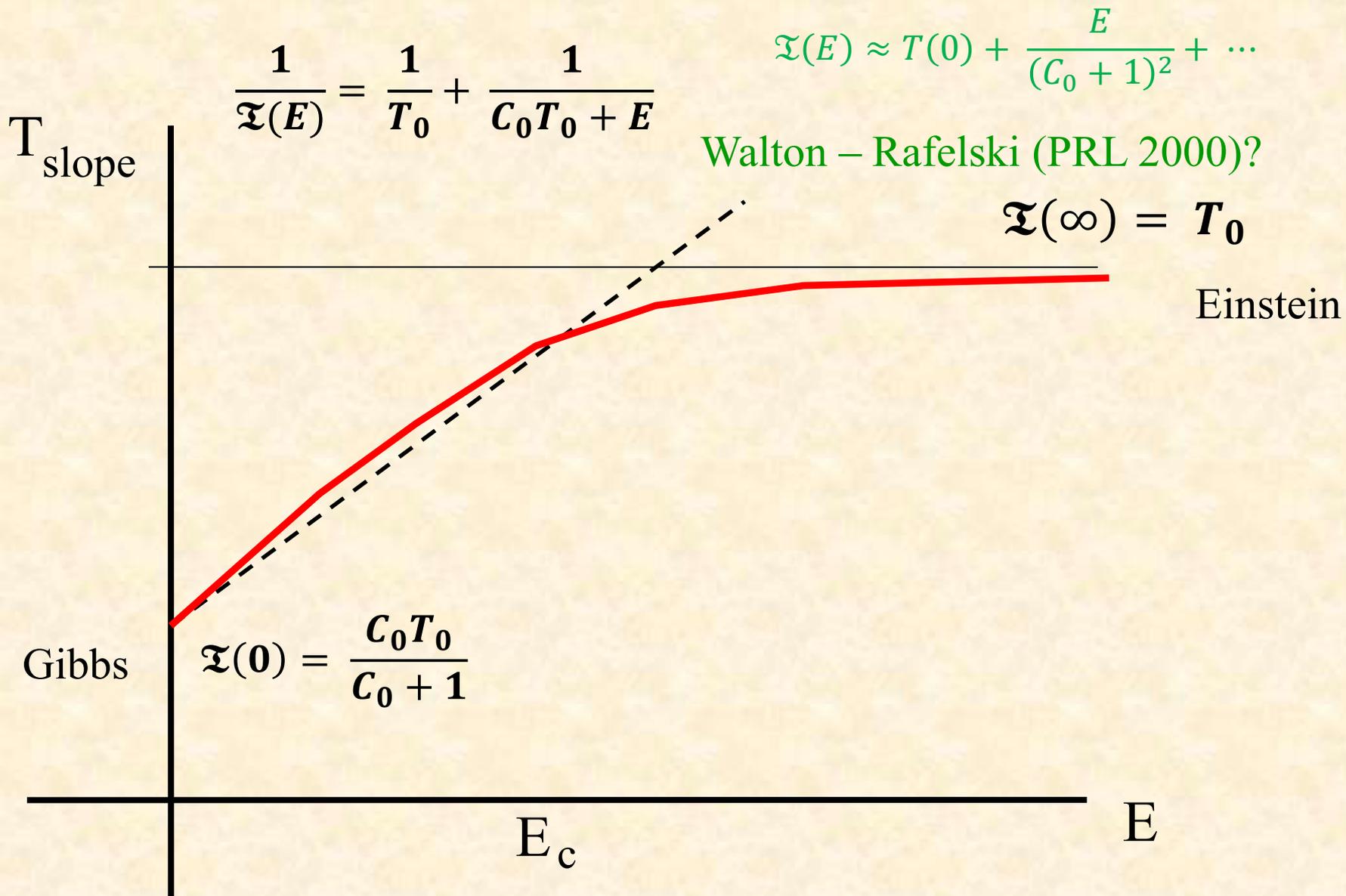
Tsallis:

$$D_{ij} (E) = \left(T + (q-1) E \right) G_{ij}$$

Field theory calculation

- polynomial interaction, one field integrated out
- Imaginary part of self-energy \rightarrow noise
- Effective Langevin eq. for soft field
- Fluctuation-dissipation: $D_{ij} = T_0 G_{ij}$
(*constant Einstein temperature*)
- G is linear in the energy: $G_{ij} = \gamma_{ij}(1 + E/T_0 C_0)$
- $f(E)$ Gaussian

$$\frac{1}{\mathfrak{Z}(E)} = \frac{1}{T_0} + \frac{1}{C_0 T_0 + E}$$



Additive and multiplicative noise

Equivalent descriptions:

TSB, AJ, PRL 94, 132302, 2005

1. Langevin

$$\dot{p} = \zeta - \gamma p$$

$$\langle \gamma \rangle = G$$

$$\langle \zeta \rangle = F$$

$$\langle \gamma \gamma \rangle_c \sim 2C \quad \langle \gamma \zeta \rangle_c \sim 2B \quad \langle \zeta \zeta \rangle_c \sim 2D$$

2. Fokker Planck

$$\frac{\partial f}{\partial t} = \frac{\partial^2}{\partial p^2} (K_2 f) - \frac{\partial}{\partial p} (K_1 f)$$

$$K_1 = F - Gp$$

$$K_2 = D - 2Bp + Cp^2$$

Exact stationary distribution:

$$f = f_0 (D/K_2)^v \exp\left(-\frac{\alpha}{\theta} \operatorname{atan}\left(\frac{\theta p}{D - Bp}\right)\right)$$

with $v = 1 + G/2C$ *power*

$$\alpha = GB/C - F \quad \textit{exponent}$$

$$\theta^2 = DC - B^2 \quad \textit{(small or large) parameter}$$

For $F = 0$ characteristic scale: $p_c^2 = D/C$.

Exact stationary distribution for $F = 0$, $B = 0$:

$$f = f_0 \left(1 + \frac{C p^2}{D} \right)^{-(1+G/2C)}$$

With $E = p^2 / 2m$ this is a Tsallis distribution!

$$f = f_0 \left(1 + (q-1) \frac{E}{T} \right)^{\frac{q}{1-q}}$$

● Tsallis index:

● Temperature:

$$q = 1 + 2C / G$$

$$T = D / mG$$

Limits of the Tsallis distribution:

$p \ll p_c$: Gauss

$$f \sim \exp(-Gp^2/2D)$$

$p \gg p_c$: Power-law

$$f \sim (p / p_c)^{-2\nu}$$

Energy distribution limits:

$$E \ll E_c: \quad f \sim \exp(-E/T)$$

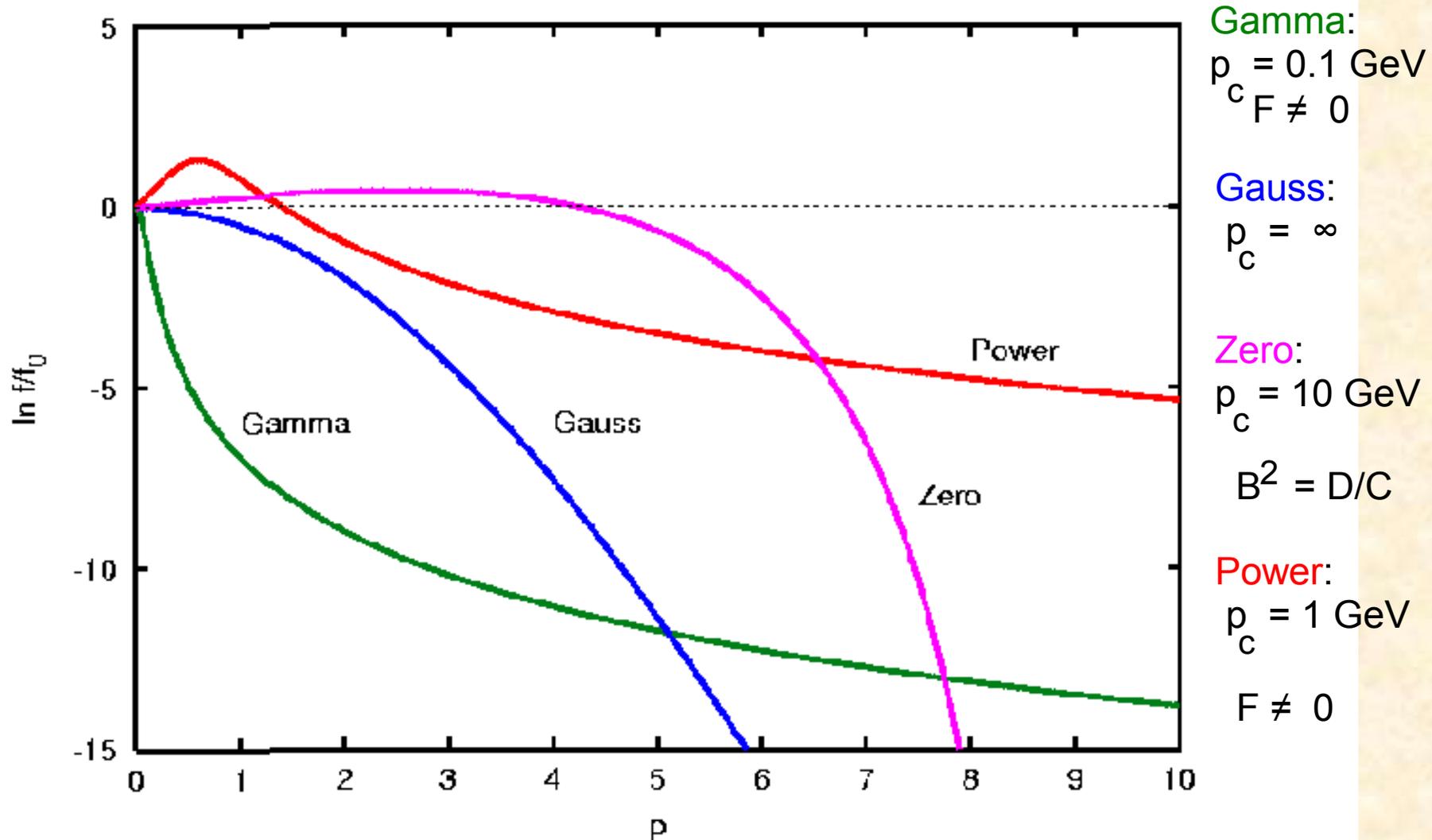
$$E \gg E_c: \quad f \sim (E/E_c)^{-\nu}$$

Relation between **slope**, **inflection** and **power** !!

$$\nu = 1 + E_c / T$$

Stationary distributions

For $F=0$, $B=0$ the Tsallis distribution is the exact stationary solution

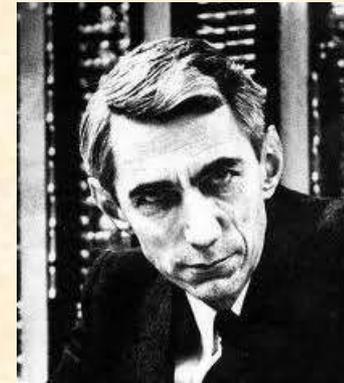


MODERN PROBLEMS

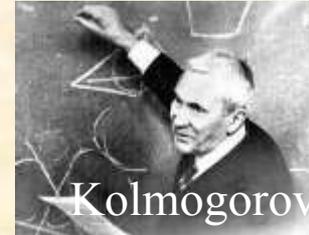
- **Information**



Neumann

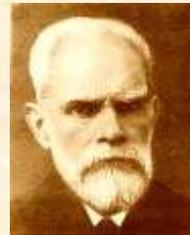


Shannon



Kolmogorov

- **Chaos**



Lyapunov



Rényi

- **Phase Transition**



- **Scaling**



Feynman



Wilson

- **Anomalous Fluctuations**



Lévy



Tsallis

Statistics, entropy, temperature

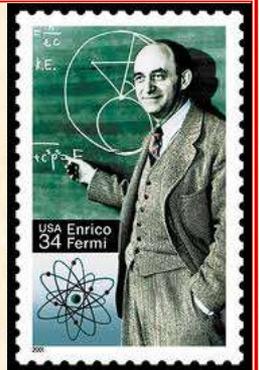
Fermi distribution (Bernoulli, Poisson)

Bose distribution (negative binomial)

Superstatistics: distribution of distributions

Fermi distribution: N, K-N

$$W = \frac{K!}{N!(K-N)!} = \binom{K}{N}$$



$$\Omega = k_B \ln W - \beta \hbar \omega (N + 1/2) + \beta \mu N = \max$$

$$\Omega_{N+1} \leq \Omega_N \quad \Omega_{N-1} \leq \Omega_N$$

Fermi distribution: $N, K-N$

Notation: $\ln x = \beta (\hbar\omega - \mu)$, $k_B = 1$.

$$\Omega_{N+1} - \Omega_N = \ln \frac{K-N}{N+1} - \ln x$$

$$\Omega_N - \Omega_{N-1} = \ln \frac{K-N+1}{N} - \ln x$$

$$\frac{K-N}{N+1} \leq x \leq \frac{K-N+1}{N}$$



Fermi distribution: N, K-N

$$\text{fix : } \bar{f} = N / K, \quad K \rightarrow \infty$$

$$\frac{1 - \bar{f}}{\bar{f} + 1/K} \leq x \leq \frac{1 - \bar{f} + 1/K}{\bar{f}}$$

$$\frac{1}{1+x} - \frac{1}{K} \frac{x}{1+x} \leq \bar{f} \leq \frac{1}{1+x} + \frac{1}{K} \frac{1}{1+x}$$

Fermi distribution: N, K-N

$$\bar{f} \rightarrow w_{\text{Fermi}}(x) = \frac{1}{1+x}$$

$$w_{\text{Fermi}} = \frac{1}{e^{\beta(\hbar\omega - \mu)} + 1}$$

Fermi distribution in a subsystem

$$P_{n,k} = \frac{\binom{k}{n} \binom{K-k}{N-n}}{\binom{K}{N}}$$

Fermi distribution in small subsystems

$$k \ll K, \quad \Rightarrow \quad n \ll N, \quad (N - n)! \approx N! N^{-n}$$

$$P_{n,k} \approx \frac{\binom{k}{n} K! K^{-k}}{\binom{K}{N} N! N^{-n} (K - N)! (K - N)^{-(k-n)}}$$

Fermi distribution in small subsystems = Bernoulli distribution

$$P_{n,k} \approx \binom{k}{n} \left(\frac{K-N}{K} \right)^k \left(\frac{N}{K-N} \right)^n$$

$$P_{n,k} \approx \binom{k}{n} \bar{f}^n (1 - \bar{f})^{k-n}$$

A story of false coins



Bose distribution in a sunsystem

$$P_{n,k} = \frac{\binom{k+n}{n} \binom{K-k+N-n}{N-n}}{\binom{K+N+1}{N}}$$

k levels and n excitations mixed arbitrarily...



Bose distribution in small subsystems

$$P_{n,k} \approx \binom{k+n}{n} \bar{f}^n (1 + \bar{f})^{-k-1-n}$$

$$\langle n \rangle = (k+1)\bar{f}$$

$$\bar{f} = w_{\text{Bose}} = \frac{1}{e^{\beta(\hbar\omega - \mu)} - 1}$$



Negative binomial distribution (NBD)

$$\binom{k+n}{n} = (-1)^n \binom{-k-1}{n}$$

$$P_{n,k} \approx \binom{-k-1}{n} (-\bar{f})^n (1 + \bar{f})^{-k-1-n}$$

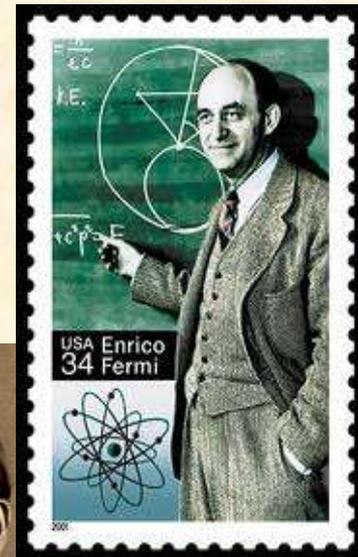
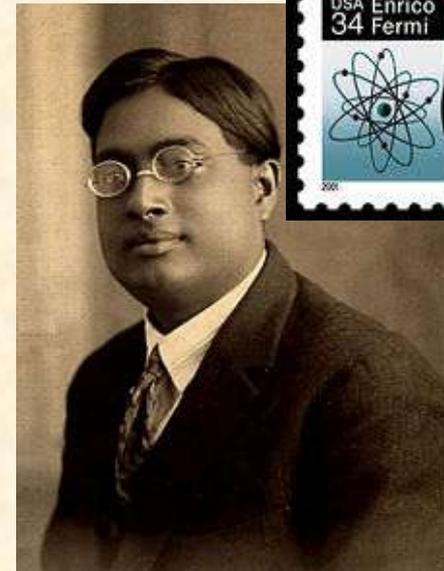


Fermi – Bose transformation: statistical supersymmetry

$$B_{n,k}(\bar{f}) = F_{n,-k-1}(-\bar{f})$$

$$F_{n,k}(\bar{f}) = B_{n,-k-1}(-\bar{f})$$

invariant $k(k+1)$



Rare events: Poisson distribution

$$P_n = \frac{a^n}{n!} e^{-a}$$

$$\langle n \rangle = a$$

$$\delta n^2 = a$$



Rare event Bernoulli: Poisson

$$n \ll k \quad \binom{k}{n} \approx \frac{k^n}{n!}$$

$$P_n \approx (1 - \bar{f})^k \frac{1}{n!} \left(\frac{k\bar{f}}{1 - \bar{f}} \right)^n$$

$$P_n = C_k(x) \frac{1}{n!} \left(ke^{-x} \right)^n$$



Rare event NBD: Poisson

$$n \ll k \quad \binom{k+n}{n} \approx \frac{k^n}{n!}$$

$$P_n \approx (1 + \bar{f})^{-k-1} \frac{1}{n!} \left(\frac{k\bar{f}}{1 + \bar{f}} \right)^n$$

$$P_n = C_k(x) \frac{1}{n!} (ke^{-x})^n$$



CORRECTIONS

- **Finite Size Effects**
- **Near-Equilibrium Fluctuations**
- **Scaling Fluctuations**
- **Superstatistics**

NBD = Euler \circ Poisson

$$\int_0^{\infty} x^N e^{-ax} dx = \frac{N!}{a^{N+1}}$$

$$P_{n,k} = \binom{k+n}{n} \bar{f}^n (1+\bar{f})^{-k-1-n} =$$

$$\frac{\bar{f}^n}{k!n!} \int_0^{\infty} x^{k+n} e^{-(1+\bar{f})x} dx$$



NBD = Euler \circ Poisson

$$P_{n,k} = \int_0^{\infty} \frac{(x \bar{f})^n}{n!} e^{-\bar{f}x} \cdot \frac{x^k}{k!} e^{-x} dx$$



Poisson in k , Euler-Gamma in x

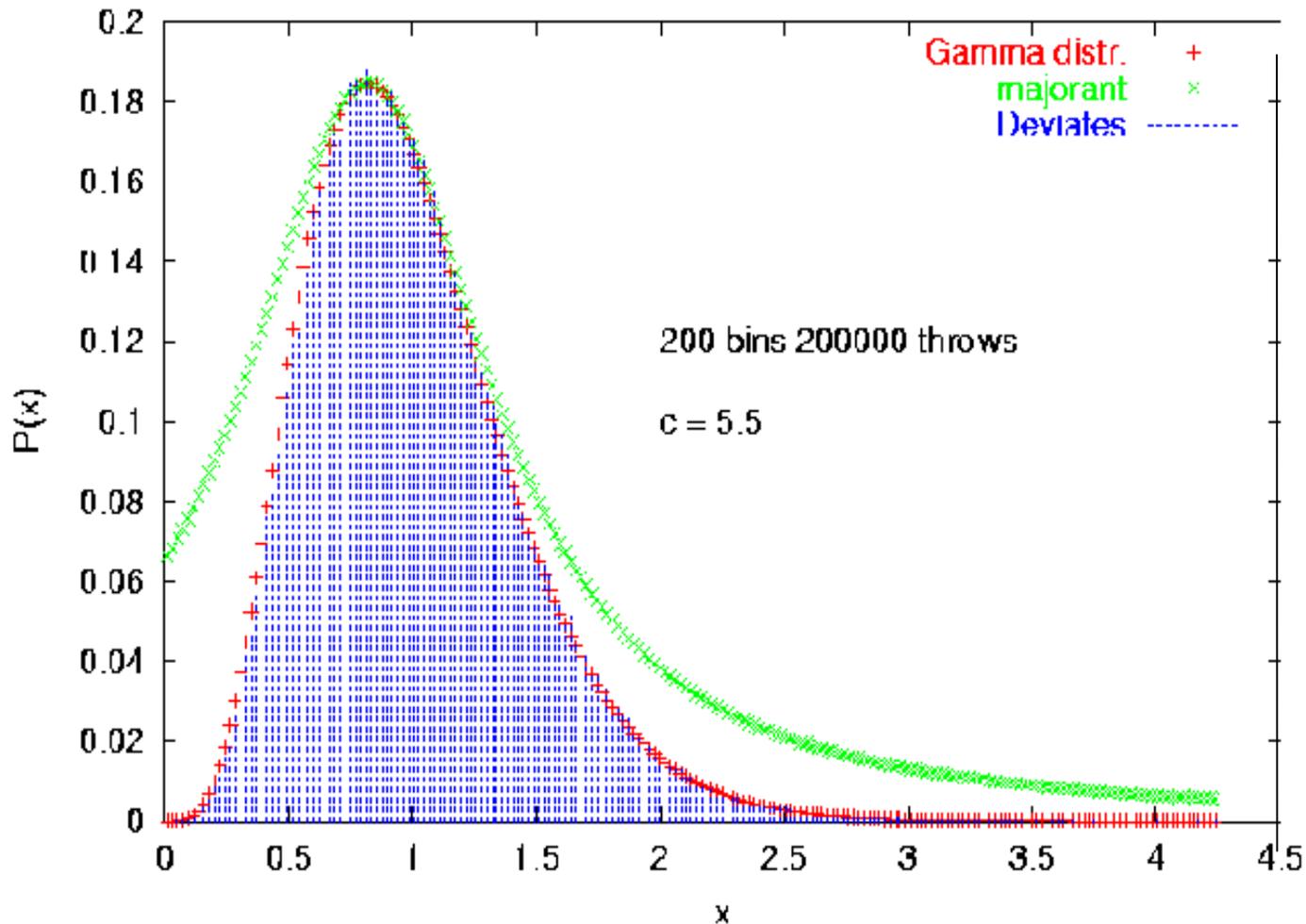


Superstatistics

Euler - Gamma distribution

max: $1 - 1/c$, mean: 1, spread: $1 / \sqrt{c}$

Gamma distributed random deviates



Homework problems

1. Regard the following distributions:

- Bernoulli ($n, k; f$)
- NBD ($n, k; f$)
- Poisson ($n, k; f$)

Questions:

- Check the norm
- n expectation value, squared variance
- characteristic function (expectation value of $\exp(bn)$)

Homework problems

1. How is the supertransformation for finite subsystems inside finite systems ?

$$B(n; k | N; K) \leftrightarrow F(n; k | N; K)$$

What is the expectation value of $f(a+x)$, if x is distributed as

a) Gauss

b) Euler-Gamma ?



Is acceleration a heat container?